Advanced quantum information: entanglement and nonlocality

Alexander Streltsov

7th class April 20, 2022

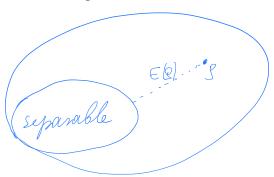
Advanced quantum information (6th class)

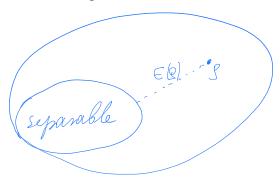
- Every Wednesday 15:15 17:00
- Literature:
 - Nielsen and Chuang, Quantum Computation and Quantum Information, Cambridge University Press (2012)
 - Horodecki et al., Quantum entanglement, Rev. Mod. Phys. 81, 865 (2009)
- Part 1 lecture materials: http://qot.cent.uw.edu.pl/teaching/
- Part 2 lecture materials: http://jkaniewski.fuw.edu.pl/?q=teaching

Quantification of entanglement
 Distance-based entanglement measures
 Negativity
 Distillable entanglement and entanglement cost

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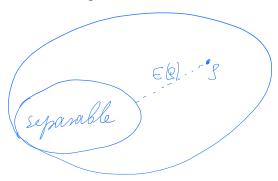




For a distance function $D(\rho, \sigma)$ define

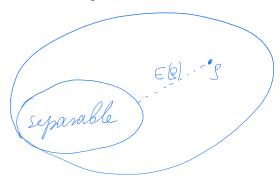
$$E(\rho) = \inf_{\sigma \in \mathcal{S}} D(\rho, \sigma)$$

with infimum over separable states ${\cal S}$



E is an entanglement measure if:

1 $D(\rho, \sigma) \ge 0$ with equality for $\rho = \sigma$

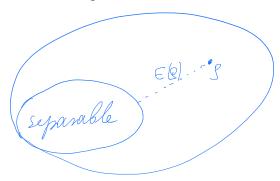


E is an entanglement measure if:

- **1** $D(\rho, \sigma) \ge 0$ with equality for $\rho = \sigma$
- **2** *D* fulfills the data-processing inequality:

$$D(\Lambda[\rho], \Lambda[\sigma]) \le D(\rho, \sigma)$$

for any quantum operation Λ



Exercise: prove that

$$E(\rho) = \inf_{\sigma \in \mathcal{S}} D(\rho, \sigma)$$

is an entanglement measure

Proof that $E(\rho) = \inf_{\sigma \in \mathcal{S}} D(\rho, \sigma)$ does not increase under LOCC:

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- We have

$$E\left(\Lambda_{LOCC}[\rho]\right) = \min_{\mu \in \mathcal{S}} D\left(\Lambda_{LOCC}[\rho], \mu\right) \le D\left(\Lambda_{LOCC}[\rho], \Lambda_{LOCC}[\sigma]\right)$$
$$\le D(\rho, \sigma) = E(\rho)$$

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• Proof holds also if Λ_{LOCC} is replaced by separable operations

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- · Upper bound on distillable entanglement
- For pure states: $E_r(|\psi\rangle^{AB}) = S(\rho^A)$
- For mixed states: hard to compute in general

Examples for distances fulfilling $D(\Lambda[\rho], \Lambda[\sigma]) \leq D(\rho, \sigma)$:

Bures distance

$$D_b(\rho,\sigma) = \sqrt{2-2F(\rho,\sigma)}$$

with fidelity
$$F(\rho, \sigma) = \operatorname{Tr} \sqrt{\sqrt{\rho}\sigma\sqrt{\rho}}$$

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Trace distance

$$D_t(\rho,\sigma) = \frac{1}{2} \|\rho - \sigma\|_1$$

with the trace norm $||M||_1 = \text{Tr } \sqrt{M^{\dagger}M}$

Quantification of entanglement

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Negativity

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• It holds that $E_n(\rho^{AB}) \ge 0$, and $E_n(\rho^{AB}) = 0$ if ρ^{AB} has non-negative partial transpose

Theorem 6.2. Negativity does not increase under LOCC:

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Proof.

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$$\sum_{i} A_{i}^{\dagger} A_{i} \otimes B_{i}^{\dagger} B_{i} = \mathbb{1}_{AB}$$

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 $\Rightarrow A_i \otimes B_i^*$ are also valid Kraus operators

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Proof.

• Partial transpose of $\Lambda_{LOCC}[\rho^{AB}]$:

$$\left(\Lambda_{\text{LOCC}} \left[\rho^{AB} \right] \right)^{T_B} = \left(\sum_i A_i \otimes B_i \rho^{AB} A_i^{\dagger} \otimes B_i^{\dagger} \right)^{T_B}$$

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Taking the trace norm gives

$$\left\| \left(\Lambda_{\text{LOCC}} \left[\rho^{AB} \right] \right)^{T_B} \right\|_{1} = \left\| \sum_{i} A_i \otimes B_i^* \rho^{T_B} A_i^{\dagger} \otimes B_i^{T} \right\|_{1} = \left\| \tilde{\Lambda} \left[\rho^{T_B} \right] \right\|_{1}$$

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• Trace norm monotonic under quantum operations:

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In summary:

$$\left\| \left(\Lambda_{\text{LOCC}} \left[\rho^{AB} \right] \right)^{T_B} \right\|_1 \le \left\| \rho^{T_B} \right\|_1$$

Monotonicity of E_n under LOCC

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Recall definition of negativity:

$$E_n(\rho^{AB}) = \frac{||\rho^{T_B}||_1 - 1}{2}$$

Q.E.D.

Negativity

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Negativity is strongly monotonic:

$$\sum_{i} q_{i} E_{n} \left(\sigma_{i}^{AB} \right) \leq E_{n} \left(\rho^{AB} \right)$$

for any states σ_i^{AB} and probabilities q_i obtainable from ρ^{AB} via LOCC

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Negativity is convex:

$$E_n\left(\sum_i p_i \rho_i^{AB}\right) \leq \sum_i p_i E_n\left(\rho_i^{AB}\right)$$

Outline

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2 Entanglement monogamy

• Distillable entanglement: singlet rate obtainable from a quantum state ρ via LOCC in the asymptotic limit

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- · Explicit formula:

$$E_d(\rho) = \sup \left\{ r : \lim_{n \to \infty} \left(\inf_{\Lambda} \left\| \Lambda \left[\rho^{\otimes n} \right] - |\Phi^+\rangle \langle \Phi^+|^{\otimes \lfloor rn \rfloor} \right\|_1 \right) = 0 \right\}$$

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- Entanglement cost: singlet rate required to create a state ρ via LOCC in the asymptotic limit
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$$E_c(\rho) = \inf \left\{ r : \lim_{n \to \infty} \left(\inf_{\Lambda} \left\| \rho^{\otimes n} - \Lambda \left[|\Phi^+\rangle \langle \Phi^+|^{\otimes \lfloor rn \rfloor} \right] \right\|_1 \right) = 0 \right\}$$

 E_d and E_c are special cases of asymptotic state-conversion rates

$$R(\rho \to \sigma) = \sup \left\{ r : \lim_{n \to \infty} \left(\inf_{\Lambda} \left\| \Lambda \left[\rho^{\otimes n} \right] - \sigma^{\otimes \lfloor rn \rfloor} \right\|_{1} \right) = 0 \right\}$$

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• It holds

$$E_d(\rho) = R(\rho \to |\Phi^+\rangle\langle\Phi^+|), \quad E_c(\rho) = \left[R(|\Phi^+\rangle\langle\Phi^+|\to\rho)\right]^{-1}$$

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• For pure states holds

$$R(\ket{\psi}
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ho_{\psi})}{S(
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where ρ_{ψ} is the reduced state of $|\psi\rangle$

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$$E_d(\rho^{AB}) \le E_c(\rho^{AB}) \le E_f(\rho^{AB})$$

Bounds on E_d and E_c :

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Application: consider maximally correlated state

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$$\rho_{\rm mc}^{\sf AB} = \sum_{i,j} \alpha_{ij} \, |ii\rangle\langle jj|$$

For $\sigma_{\text{sep}}^{AB} = \sum_{i} \alpha_{ii} |ii\rangle\langle ii|$ it holds

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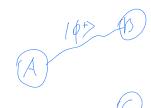
$$\Rightarrow \textit{E}_{\textit{d}}(\rho_{\textit{mc}}^{\textit{AB}}) = \textit{S}(\rho_{\textit{mc}}^{\textit{A}}) - \textit{S}(\rho_{\textit{mc}}^{\textit{AB}})$$

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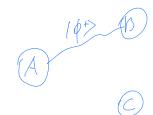
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Solution:

• Consider total state ρ^{ABC}

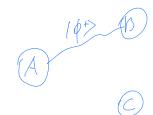
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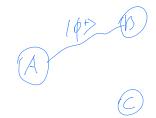
- Consider total state ρ^{ABC}
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- ⇒ total state must be

$$\rho^{\mathit{ABC}} = |\Phi^+\rangle\langle\Phi^+|^{\mathit{AB}}\otimes\rho^{\mathit{C}}$$

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Entanglement monogamy: If two qubits *A* and *B* are maximally entangled, they cannot be entangled with another qubit *C*



Compare to classical random variables: a classical random variable *A* can be maximally correlated with *B* and *C* at the same time:

$$\rho^{ABC} = \frac{1}{2} |000\rangle\langle000| + \frac{1}{2} |111\rangle\langle111|$$

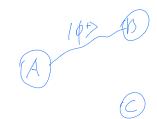
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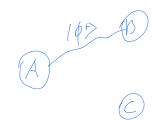


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• $C_{A:B}$ and $C_{A:C}$: concurrence of the reduced state ρ^{AB} and ρ^{AC}

•
$$C_{A:BC} = \sqrt{2(1 - \text{Tr}[(\rho^A)^2])}$$